

**Homework #2** — Phys625 — Spring 2002  
**Deadline: Thursday, February 14, 2002.**

Turn in homework in the class or put it in  
the box on the door of Phys 2314 by 10 a.m.

Web page: <http://www2.physics.umd.edu/~yakovenk/teaching/phys625.spring2002>

Victor Yakovenko, Associate Professor  
Office: Physics 2314  
Phone: (301)-405-6151  
E-mail: yakovenk@physics.umd.edu

**Do not forget to write your name and the homework number!**

Equation numbers with the period, like (3.25), refer to the equations of the textbook.

Equation numbers without period, like (5), refer to the equations of this homework.

## Fermi Liquid Theory

### 1. *Plasma oscillations*

- (a) [6 points] Coulomb interaction between electrons gives the following nonlocal-in-space correction to the energy of quasiparticle:

$$\delta\epsilon(\mathbf{r}) = \int \frac{d^3\mathbf{r}' d^3\mathbf{p}}{(2\pi)^3} \frac{e^2 \delta n(\mathbf{p}, \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} \quad (1)$$

Repeat the derivation of zero sound given in §4 using Eq. (1). Show that Eq. (4.15) acquires the following form:

$$\frac{\omega}{2v_F k} \ln \left( \frac{\omega + v_F k}{\omega - v_F k} \right) - 1 = \frac{k^2}{4\pi e^2 \nu_F}, \quad (2)$$

where  $\nu_F$  is the density of states given by Eq. (2.5).

From Eq. (2), determine the frequency  $\omega_0$  of the collective mode in the limit  $k \rightarrow 0$ . It is called the plasma frequency.

- (b) [2 points] Let us derive the frequency of plasma oscillations in a different way. Consider the continuity equation for the electron gas:

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{v}) \approx \frac{\partial \delta n}{\partial t} + n_0 \nabla \cdot \mathbf{v} = 0, \quad (3)$$

where  $n = n_0 + \delta n$ , and  $\mathbf{v}$  is the velocity of hydrodynamic flow. Let us differentiate Eq. (3) in time:

$$\frac{\partial^2 \delta n}{\partial t^2} + n_0 \nabla \cdot \frac{\partial \mathbf{v}}{\partial t} = 0. \quad (4)$$

Using Newton's equation of motion  $\partial \mathbf{v} / \partial t = -e\mathcal{E}/m$  and Poisson's equation  $\nabla \cdot \mathcal{E} = -4\pi e \delta n$ , where  $\mathcal{E}$  is electric field, transform Eq. (4) to the following form

$$\frac{\partial^2 \delta n}{\partial t^2} = -\omega_0^2 \delta n \quad (5)$$

and determine the frequency  $\omega_0$ . Eq. (5) describes local oscillations of electron density.

**2. [8 points]** *Fermi-liquid modes in 2D*

Let us consider a Fermi liquid in two-dimensional (2D) space, where the Fermi sphere is a circle. The deviation of the quasiparticle distribution function  $\nu(\theta)$  is a function of the angle  $\theta$  on the circle. Suppose the Landau interaction function is a constant:  $F(\mathbf{n} - \mathbf{n}') = F_0$ . Using the Fourier expansion  $\nu(\theta) = \sum_m \nu_m \exp(im\theta)$ , show that the kinetic equation becomes

$$\omega \nu_m = \frac{1}{2} k v_F (\tilde{\nu}_{m+1} + \tilde{\nu}_{m-1}), \quad \text{where} \quad \tilde{\nu}_m = \begin{cases} \nu_m, & \text{when } m \neq 0, \\ (1 + F_0) \nu_0, & \text{when } m = 0. \end{cases} \quad (6)$$

Because Eq. (6) is invariant with respect to translations  $m \rightarrow (m \pm 1)$ , except at  $m = 0$ , we can look for its solutions in the form  $\nu_m = A^\pm \exp(i\alpha m)$  for  $m > 0$  and  $m < 0$ , and then match them at  $m = 0$ .

Show that for any real  $0 < \alpha < \pi$ , there exist two such solutions describing quasiparticles with  $\omega = v_F k \cos \alpha$  and  $\theta = \pm \alpha$ . [By parity, the two solutions can be written as  $\nu_m = \sin(\alpha m)$  and as  $\nu_m = \cos(\alpha m \pm \lambda)$  for  $m > 0$  and  $m < 0$ , and the parameter  $\lambda$  needs to be determined.]

Besides, for  $F_0 > 0$ , there exist one solution with a complex  $\alpha$ , which decays exponentially at  $m \rightarrow \pm \infty$ . It corresponds to the zero-sound collective mode. Find its velocity  $\omega/k$  and angular dependence  $\nu(\theta)$ . Show that the frequency of this solution lies above the quasiparticle continuum  $\omega(k) > v_F k$ .

[Notice a mathematical analogy between this problem and the problem of a 1D quantum particle subject to the  $\delta(x)$  potential.]

**3. [6 points]** *Electron-hole continuum*

Let us introduce the operator  $\hat{c}_{\mathbf{p},\mathbf{k}}^+$  that creates a hole with momentum  $\mathbf{p}$  and an electron with momentum  $\mathbf{p} + \mathbf{k}$  out of the Fermi sea:

$$\hat{c}_{\mathbf{p},\mathbf{k}}^+ = \hat{a}_{\mathbf{p}+\mathbf{k}}^+ \hat{a}_{\mathbf{p}}, \quad \text{where} \quad |\mathbf{p}| < p_F \quad \text{and} \quad |\mathbf{p} + \mathbf{k}| > p_F. \quad (7)$$

The energy of such an excitation is

$$E_{\mathbf{p},\mathbf{k}} = \epsilon_{\mathbf{p}+\mathbf{k}} - \epsilon_{\mathbf{p}} = \frac{|\mathbf{p} + \mathbf{k}|^2 - |\mathbf{p}|^2}{2m}. \quad (8)$$

With the restriction (7) on  $\mathbf{p}$  and  $\mathbf{k}$ , outline the area covered by these excitation on a plot of  $E(k)$  for all permitted values of  $p$ . This is the so-called electron-hole continuum of excitations. On the same plot, also sketch the dispersions  $\omega(k)$  of the zero-sound and plasma-oscillations collective modes.